THE VIBRATIONAL SPECTRA OF VINYLMETHYLHALOSILANES $H_2C=CH(CH_3)_{3-n}SiX_n$ (n = 1, 2; X = F, Cl)* **

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The infrared spectra (4000–200 cm⁻¹) of gaseous and liquid or dissolved samples of vinyldimethylfluorosilane (C_4H_9SiF , I), vinylmethyldifluorosilane ($C_3H_6SiF_2$, II), and of the chloroanalogs (C_4H_9SiCI , III and $C_3H_6SiCI_2$, IV) were measured together with the laser Raman spectra of the liquids (1700–100 cm⁻¹). The 70 eV mass spectra of I and II were also recorded. The assignment was carried out only for bands above 500 cm⁻¹; it indicated presence of rotamers.

Existence of pairs of stable rotamers with the *skew* (antiperiplanar) and *cis* (synplanar) conformations has been already predicted¹ for all molecules which represent the combinations of a *sp*³-hybridized asymmetric top (*e.g.* secondary, tertiary, and quarternary carbon atoms) and a *sp*²-hybridized frame (*e.g.* vinyl group).

As far as we know, no attempt to recognize rotamers in silicon series had been made though many substituted vinylsilanes were synthetized. For instance, the *IR* spectra of both liquid *III* and *IV* had been published² in figures ($3200-400 \text{ cm}^{-1}$), however, the assignment³ was not carried out bellow 900 cm⁻¹. The only exception had been identification⁴ of absorptions by the Si—Cl stretching vibration (477 for *III*, 462 and 562 cm⁻¹ for *IV*). The tabulated Raman spectrum without polarizations had been quoted⁵ only for *III* but without warranty for its purity.

The aim of this work was to re-examine the published spectra of III, IV and to interpret our spectra of I-IV especially in the 900-500 cm⁻¹ region.

EXPERIMENTAL

I and II were prepared by reaction of III and IV with SbF₃ respectively, their physical constants and detailed synthesis had been published⁶ elsewhere. Purity of I and II was checked by gas-liquid chromatography and mass spectrometry. According to these methods, the content of the parent chloro compounds was less than 0·1 mol. %. The mass spectra were measured on a LKB spectrometer, model GC-MS 9000 at 70 and 13 eV, the temperature of inlet system being 150°C, that of ion source 270°C. Ten peaks at 70 eV are collected in the decreasing order of intensities (in per cent of height of the base peak).

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- *I* (m.w. = 104): m/e = 89 (100), 63 (51), 47 (31), 77(21), 104 (17), 49 (9.6), 90 (8.9), 76(5.6), 62 (4.4), 91 (3.9); $m^{++} = 44.6 (89^+ \rightarrow 63^+ + 26).$
- *II* (m.w. = 108): m/e = 93 (100), 81 (55), 108 (53), 67 (28), 47 (27), 42 (21), 27 (15), 94 (7-9), 89 (7-9), 80 (7-9), 91 (7-5); $m^{++} = 48\cdot3 (93^+ \rightarrow 67^+ + 26).$

The absorption spectra were measured on double-beam spectrometers Beckman IR-7 (700 to 200 cm⁻¹, Cst-prism/grating optics) and Zeiss (Jena) UR-20 (4000-400 cm⁻¹, prisms LiF, NaCl, KBr). Spectra of gases were taken in a 10 cm cell with KRS-5 or KBr windows. A droplet of liquid sample was introduced into the cell during flushing it with dried nitrogen. Then, both stopcocks were turned off and the sample was allowed to evaporate. *II* (b.p. 25°C) was cooled

TABLE I

Infi	rared	R	aman	- Tentative assignment
gas ^b	type(I)	liq.	I (dep.)	- Tentative assignment
1 933	asym. (w)	NM	NM	$2 \text{ x wag.} = CH_2$
1 603	PQR (m)	1 598	7 (0.4)	stretching C=C
1 409	PQR (s)	1 410	11 (0.5)	asym. def. $SiMe_2 + sciss. = CH_2$
(1 270)	sh↓ (?)	1 275	11 (0.4)	in-plane def. =-C-H
1 265	(s)			sym, def. SiMe ₂
1 2 2 0	sharp (s)			in-plane def. = C-H
(1 078)	(m)			rocking == CH ₂ ?
1013d	(s)		1	
1 007 J. a	(s)		}	twisting HC=-CH
963	PQR (s)			wagging $= CH_2$
950	sh↑ (?)			?
895	PQR (vs)		1	
(878)	(vs)		}	stretching Si—F
850	PQR (vs)		۱	and in the SiMe
799	PQR (vs)	794	4 (1∙0)∫	rocking SiMe ₂
774	(vs)	768	1 - 2(p ?)	
709	PQR (s)		}	asym. stretch. SiC ₃
(700)	sh↑ (?)	699	7 (0·7) ^J	
(608)	(w)	610	100 (0.32)	sym. stretch. SiC ₃
522	(s)			o-o-p def. ==CH
395	PQR (s)	396	7(p?)	
369	sh↑ (?)	369} ^d	11 (0.5)	
275	PQR (s)	271)	14 (dp?)	skeletal def.
		258}d	11 (p ?)	
NM	NM	193	33 (0·87)J	

Assignment of the Strongest Bands in Vinyldimethylfluorosilane I, $(cm^{-1})^a$

^a Meaning of abbreviations: NM = not measured, d = doublet, sh \uparrow or sh \downarrow = shoulder on the preceding or following band, o-o-p = out-of-plane, Me = methyl; ^b wave number in parentheses is for liquid or solution in CS₂. The band center is given for *PQR*-bands.

in a glass baloon before introduction. This procedure usually offered a high concentration at which the weakest bands were recorded. Pressure in the cell was then gradually decreased by evacuation so that the absolute pressures were not known. After all measurements, moist air was introduced into the cell and the bands of hydrolytic products were identified. IR spectra of liquids were measured as capillary films between KBr plates, for *II*, moreover, a cold solution in CS_2 was investigated in an incompletely filled 0.1 cm KBr cell. The instruments were calibrated with water vapour, CO_2 , and 10 μ polystyrene film, spectra of *III* and *IV* also with gaseous HCI formed by their hydrolysis. The Raman spectra with polarizations were measured with He—Ne laser excitation at room temperature in the 1700–100 cm⁻¹ region, for *II* in a cooled cell at $-55^{\circ}C$. Accuracy of readings was ± 2 cm⁻¹.

RESULTS

The survey spectra of I and II are in Figs 1, 2, the assignments for I-IV are in Tables I-IV. The IR bands of hydrolytic product in gaseous and liquid I-IV are collected in Table V.

The strongest Raman lines (polarized) of vinylfluorosilanes are 610 (0.32) for *I*, 641 (0.40) for *II*, and 692 cm⁻¹ (0.28) for vinyltrifluorosilane⁷. The same trend we observed for the series of ethylmethylfluorosilanes $C_2H_5(CH_3)_{3-n}SiF_n$ (n = 1, 2, 3)⁶:

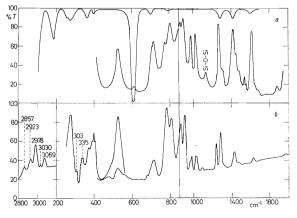
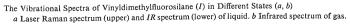


FIG. 1



600 (0.36), 625 (0.33), and 705 cm⁻¹ (0.34) respectively. It follows from this fact that no specific band has to be expected for the stretching vibration of the Si—C= bond because all $(Si-C)_{4-n}$ bonds (n = 1, 2) are vibrationally coupled. The mentioned polarized Raman lines must correspond to the totally symmetrical mode and the remaining out-of-phase Si—C stretching frequencies should be found above the respective in-phase fundamental and bellow 900 cm⁻¹.

The rocking vibration of methyl groups attached to silicon atom $(850-800 \text{ cm}^{-1})^8$ and possible existence of rotamers, along with the above mentioned fact, can cause a rather complex pattern of IR spectra in the $900-600 \text{ cm}^{-1}$ region. This holds for all the compounds studied except gaseous *I* which has a minimal number of IR bands in this region.

Vinyldimethylfluorosilane (I). For this gas, the absorptions by the vinyl group (C=C stretching, =CH₂ scissoring and wagging, twisting HC=CH, o-o-p and in-plane =C-H deformation) were identified without difficulty, the =CH₂ wagging frequency also with the aid of its first overtone at 1933 cm⁻¹. The Si-F stretching mode was assigned to the strong IR band at 895 cm⁻¹ (cf.^{7,9}), the equally intense

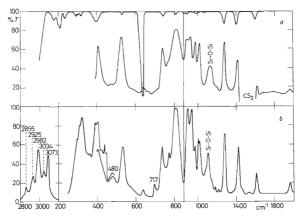


FIG. 2

The Vibrational Spectra of Vinylmethyldifluorosilane (II) in Different States (a, b)

a Laser Raman spectrum of liquid at -55° C (upper) and *IR* spectrum of a cold solution in CS₂ (lower). *b* Infrared spectrum of gas.

bands near 850 and 800 cm⁻¹ to the rocking modes of methyls in the Si(CH₃)₂ group. The band corresponding to the symmetrical deformation of this grouping (near 1260, cf.⁸) could be splitted into a doublet depending on molecular symmetry. The published data⁵ for gaseous (CH₃)₂SiFBr (C_s symmetry) show that the splitting of the 1265 cm⁻¹ band does not exceed 10 cm⁻¹, the same result we obtained with gaseous (CH₃)₂HSiCl (cf.¹⁰). In the case of gaseous analog *III* no splitting was observed at all and, therefore, we ascribed the unsplitted band at 1265 cm⁻¹ to the mentioned vibration.

TABLE	п
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Infi	rared		R	aman	Tratation and another
gas ^b	type	(I)	liq.	I (dep.)	Tentative assignment
1 957	asym.	(w)	NM	NM	$2 \times \text{wag.} = CH_2$
1 605	PQR?	(m)	1 604	10 (0.43)	stretching C=C
1 413	PQR?	(s)	1 413	15 (0.60)	asym. def. SiMe + sciss. ==CH ₂
1 273	PQR ?	(s)			sym. def. SiMe
1 218		(w)	1 280	13 (0.50)	in-plane def. ==CH
1 013 1 007}d		(s) (s)		}	twisting HC=CH
970	PQR	(s)			wagging $= CH_2$
933	PQR	(vs)			asym. stretch. SiF ₂
894		(vs)			sym. stretching SiF ₂
810	PQR	(vs)	812	4 (0.86)	rocking SiMe
770	sh 🏌	(?)		}	asym. stretch. SiC ₂
734	PQR	(s)	745	7 (0.54)∫	. 2
644		(m)	641	100 (0.40)	sym. stretch. SiC_2
528		(s)	532	5 (1.0)	o-o-p def. ≕C—H
434	PQR	(m)	200	12 (0.82)	
$\frac{404}{202}$		(s)	399	12 (0.83)	
392}d		(s) (s)	323	7 (1.0)	
321	sh↑	(s) (?)	302)	6 (dp?)	
297	sh ↑	(?)	292 d	8 (p?)	skeletal def.
NM	NM	(.)	237 sh ⊥	9 (0.75)	
NM	NM		209)	19 (1.0)	
NM	NM		201^{201} d	17 (p?)	
NM	NM		177	11 (1.0)	
NM	NM		51	vs (NM)	
NM	NM		35	vs (NM)	

Assignment of the Strongest Bands in Vinylmethyldifluorosilane II, $(cm^{-1})^a$

^{a,b} For meanings of symbols, see Table I.

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The two out-of-phase stretching modes of SiC₃ grouping can then be identified as the two strong IR absorptions at 774 and 709 cm⁻¹. Having not calculated the potential energy distribution, we did not try to assign the IR bands of gas below 500 cm⁻¹. Under these conditions all assumed group frequencies are identified with the exception of the =CH₂ rocking mode which is weak both in the IR and Raman spectra and masked by the "hydrolytic" band in liquid.

The only case when maxima of P.R-envelopes were resolved, is that for the C=C stretching band with the Q-branch at 1603 cm⁻¹. The P-R separation was 13 cm⁻¹, the maxima being 1596 and 1609 cm⁻¹, respectively. With I, on going from gas to

Infrar	ed		Raman		m
gas ^b	(I)	liq." (I)	liq. ^d	I $(dep.)^d$	- Tentative assigment
1 913	(vw)		NM	NM	$2 \times wag. = CH_2$
1 599	(w)	1 596 (7)	1 599	16 (0.32)	stretching C=C
1 410	(s)	1 406 (8)	1 410	21 (0.67)	asym. def. SiMe ₂
1 383	(m)		1 379	2(p?)	scissoring $= CH_2$
		1 270 (7)	1 275	24 (0.34)	in-plane def. = C-H
1 260	(vs)				sym. def. SiMe ₂
1 010	(s)	1 009 (2)	1 012	4 (1.0)	twisting HC=CH
958	(s)				wagging $= CH_2$
844	(vs)		856	2(?)	rocking SiMe ₂
			837	4(?) 🐧	-
(820)	(vs)				asym. stretch. SiC ₃
797	(vs)	791 (4)	794	7 (?)	rocking SiMe ₂
703	(s)	715 (1)	716	15 (0·7)	
		698 (2)	702} ^d	12 (1∙0) ∫	asym. stretch. SiC_3
635	(m)	628 (7)	634	ן (0.32)	
630} ^d				}	sym. stretch. SiC ₃
			595	18 (0.2)	? impurity SiOSi
526	(s)	518 (4)	527	(0·45) 17	o-o-p def. ==CH + str. SiCl
488	(m)	571 (8)	474	91 (0·33)∫	stretching Si-Cl
		368 (1)	366 sh ↓	16 (0.5)	
351	(s)	349 (3)	354	36 (0.3)	
240	(vs)			1	skeletal def.
NM		206 (2)	210	26 (1.0)	
NM			182	65 (1.0)	
NM			105	100 (?)	

^{a,b} For meanings of symbols, see Table I; ^c data of Engelhardt⁵ with Hg-excitation (4358 Å), intensity scale 0-10; ^d our data with laser. Additional Raman bands: 436(4, ?), 552(6, ?), 611 (8, ?), 658 (8.5, ?), 767 (3.5, ?), 806 (6, ?), 1000 (4, dp) cm⁻¹.

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TABLE III

TABLE IV

liquid, there were three remarkable changes in the IR spectra. The band at 1220 cm^{-1} was shifted to about 1270 cm^{-1} and it lowered its intensity, the absorption by the Si—F stretching (895 cm⁻¹ in gas) appeared for liquid at 878 cm⁻¹, and the bands in the 780-700 cm⁻¹ region interchanged their intensities.

Vinylmethyldifluorosilane (II). On the basis of the preceding results we assigned the bands of II. A general feature of Raman spectra of both I and II is very weak intensity of lines for the Si—F stretching vibration in comparison with corresponding absorptions in the infrared. Also noticeable is the pronounced intensity of Raman lines for the in-phase SiC₂ stretching vibration caused by presence of fluorine atoms in the molecules (cf. vinylchlorosilanes in Table III and IV).

Vinylchlorosilanes III and IV. All IR bands of both gases III and IV had simple

Infrare	d	Ran	nan ^c	Tentative assignment	
gas ^b	(I)	liq.	I (dep.)		
1 953	w	NM	NM	$2 \text{ x wag.} = CH_2$	
1 610	m	1 600	9 (0.30)	stretching C==C	
1 410	s			asym. def. SiMe	
(1 382)	m	1 405	12 (0.20)	scissoring $= CH_2$	
		1 271	15 (0.33)	in-plane def. ==CH	
1 267	s			sym. def. Si—Me	
1 004	s			twisting HC==CH	
973	s			wagging $= CH_2$	
804	vs			rocking SiMe	
(790)	vs		۱	asym. stretch. SiC_2	
754	vs	748	9 (0∙45)∫		
678	vs	679	18 (0.36)	sym. stretch. SiC ₂	
613	m	597	3(p?)}	asym. stretch. SiCl ₂	
573	vs		ſ	· ·	
524	s	521	11 (0.36)	o-o-p def. ==C-H	
483}d	s	479 sh ↓	35 (0.26)	sym. stretch. SiCl ₂	
467	s	461	100 (0.21)		
366	w	367	28 (0.30)		
		341} ^d	27 (0.26)		
238	vs	227	34 (0.83)	skeletal def.	
NM		201	25 (0.78)		
NM		160	71 (0·82)J		

Assignment of the Strongest Bands in Vinylmethyldichlorosilane IV, $(cm^{-1})^a$

^{*a,b*} For meanings of symbols, see Table I. ^{*c*} Additional Raman bands: doublet 546 (7, 0.6) and 558 (8, 7), 640 (2, 7) cm⁻¹.

shape with no PQR-structure even at higher resolutions so that our spectra generally agreed with the published figures² for the liquids. The only change was disappearance of the 820 cm⁻¹ band (liquid *III*) from the spectrum of gas.

The assignment for III and IV is in agreement with the published data³ for liquids but ours is more complete. The re-assignments of bands to the o-o-p =C—H deformation are collected in Table VI. Comparing our spectra of *III* with those tabulated by Engelhardt⁵, we excluded his IR bands 1182 (vw), 1143 (m), 1120 (w), 752 (m), 562 (s) and also the Raman counterparts at 1148 (3) and 562 (2) as the bands of impurities.

Note. The two assignments^{7,9} of the fundamental frequencies of vinyltrifluorosilane had been published but with a difference for the =CH₂ wagging mode. For comparison, we measured the IR spectrum of vinyltrifluorosilane in a cooled CS₂ solution. Complete agreement with the data of Crowder and Smyrl⁹ was reached: the wagging

Compound	Ι	11	III	IV
Gas	1 120-1 100	1 102	1 075	1 100-1 080
Liquid	1 1 20 - 1 1 00	1 1 3 0	1 070-1 050	1 1 30 1 0 50

TABLE V Infrared Bands^a of Hydrolytic Products in I-IV (cm⁻¹)

^a Always broad.

TABLE VI

The Infrared Band of the =-C-H o-o-p Deformation in Gaseous H_2C =CH(CH₃)_{n-3} SiX_n (cm⁻¹)^a

X / n	1	2	3
F	522	528 524 ^d	542 ^b
Cl	526 ^c	524 ^d	520 ^e

^a Medium to weak intensity. For vinylsilanes generally¹³ in the 580-515 cm⁻¹ region; ^b ref.⁷; ^c earlier value³ 627 cm⁻¹. Now, presumably obscured by the Si-Cl band; ^d earlier value³ 602 or 677 cm⁻¹; ^e for liquid⁴. Earlier value for liquid was¹⁴ near 611 cm⁻¹. Not observed in gas¹⁴.

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frequency is that at 986 cm⁻¹, and not at 731 cm^{-1} as Durig and Hellams⁷ have proposed. In some gaseous samples of vinyltrifluorosilane we also observed a C-type band near 730 cm⁻¹ but of variable intensity, hence we ascribed it to impurities.

DISCUSSION

It is clear from the second column of Table I that nine bands have an apparent POR-structure. However, eight of these have P- and R-branches as unresolved shoulders on the prominent central part. The same holds for II (Table II). We assume that the blurred POR-structures of bands are also an indirect evidence for presence of rotamers. Their mixtures in vinylderivatives generally complicated any interpretation of vibrational spectra. For instance, Potts and Nyquist¹¹ have presented a non-linear correlation between the =CH₂ wagging frequency of vinylgroups and the Hammett $\sigma_{\rm P}$ -constant of their substituents. Using their data and our two wagging frequencies for gaseous I and II we get for $(CH_3)_2$ SiF and CH_3 SiF₂ groupings the σ_R -constants about (+0.20) and (+0.26), respectively. It is clear that the σ_{R} -values obtained must involve a kind of sterical contribution of the substituent which cannot be separated. The plot¹¹ had been constructed from thirteen compounds, five of which can represent liquid mixtures of rotamers (-OCH₃,-OC₂H₅,-CH₂Cl,-COOC₂H₅,-COCH₃). This possibility had not been known in that time. Now, however, it may serve as an explanation of the non-linearity of the plot and of "a rather bad scatter of the points for the higher values of $(\sigma_p - \sigma')$ " regardles the sign as it has been pointed out¹¹.

In the earlier papers^{6,12}, a dependence of molecular dipole moments of liquid $H_2C=CH(CH_3)_{3-n}SiX_n$ (n = 0, 1, 2, 3; X = F, Cl) on n has been presented. The experimental points in plots were not "well-behaving" and we believe that presence of rotamers is the main cause for it.

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